Unprecedented Structures for PLi₅, SLi₄, and SLi₆

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Theoretical calculations have shown that PLi₅, SLi₄, and SLi₆ challenge conventional notions of molecular structure and bonding.

There has been recent interest in the theoretical exploration of lithium chemistry. While a formal analogy may be drawn between the ns1 electronic structures of H and Li, there are of course very substantial differences in their overall chemical behaviour. Stoicheiometries and structures have been reported for CLi_{6} , ¹ OLi_{4} , ² and $SiLi_{4}$ ³ which bear no resemblance to analogous hydrogen-containing compounds. I now describe a theoretical study of some lithium derivatives of P and S. These compounds have not yet been considered in the literature, and their structures are predicted to be even more remarkable than those of their N and O counterparts. PLi5 has $C_{4\nu}$ symmetry, SLi₆ has D_{3d} symmetry, but a definitive structural prediction is not yet possible for SLi₄, which has almost degenerate $C_{2\nu}$ and $C_{3\nu}$ isomers. All three 'hyperlithiated' compounds are predicted to be thermodynamically stable both to atomisation and to loss of Li₂, by substantial margins, but are probably less stable than the elements in their standard states.

Geometry optimisations were performed using the Gaussian 86 program⁴ and the 3-21G(*) basis. Vibrational frequencies were calculated with this basis from analytical SCF second derivatives for each stationary point located. Better estimates of final molecular energies were obtained with the nonstandard 6-31G(*) basis (which does not contain d functions on Li) at theoretical levels up to MP4SDQ. The success of theoretical methods of this type in predicting molecular structures is impressive, and they give at least semiquantitatively useful estimates of reaction energies.⁵ Geometrical parameters for the various stationary points of PLi₅, SLi₄, and SLi₆ are displayed in Figure 1. Absolute energies are reported in Table 1 and relative energies are presented in Table 2.

Two stationary points were located for PLi₅. The $C_{4\nu}$ isomer (1) is a true minimum, whose lowest vibrational mode is at 56 cm⁻¹ (a₁), but the D_{3h} structure (2) has an e' imaginary bending motion at 50i cm⁻¹, which is an in-plane motion of the equatorial atoms leading to (1) *via* a $C_{2\nu}$ pathway. Isomer (1) is

slightly but consistently more stable than (2) at all levels of theory used here. All other phosphoranes studied to date, either by experimental or theoretical methods, have a trigonal



Figure 1. Optimized HF/3-21G(*) geometries of isomers of PLi₅, (1), (2); SLi₄, (3)–(7); and SLi₆, (8)–(11). Bond lengths in Å, angles in degrees.

Table 1. Energies of PLi _n and SLi _n molecules (-ator	nic units))
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		Point group	3-21G(*)// 3-21G(*)						
Species				Z.P.E.ª	RHF/		 MP3/	MP4DQ/	MP4SDQ/
PLi ₃		$C_{3\nu}$	361.300 80	11.8(0)	363.07123	363.205 85	363.21480	363.21485	363.21761
PLi _s	(1)	C_{4v}	376.136 69	18.1 (O)	377.99601	378.164 42	378.175 37	378.17288	
	(2)	D_{3h}	376.13535	17.0(2)	377.994 55	378.164 02	378.17461	378.171 76	
SLi ₂	()	D_{mh}^{sn}	410,53778	8.1 (O)	412.474 34	412.60161	412.61192	412.61228	412.613 50
SLi	(3)	C_{3y}	425,362.92	14.6(0)	427.388 82	427.54721	427.56331	427.564 59	427.56618
4	(4)	$C_{2\nu}$	425,361 62	13.9(1)	427.38738	427.54943	427.56476	427.56506	427.56674
	(5)	$C_{2\nu}$	425,36033	13.3(2)	427.38598	427.54867	427.563 85	427.563 98	427.565 58
	6	\tilde{T}_{d}	425,358.85	12.5(3)	427.38296	427.54810	427.56171	427.56042	427.56192
	$(\tilde{7})$	D_{Ab}	425,34623	12.8(2)	427.37212	427.540 52	427.55524		
SLis	(8)	D_{2d}	440,151 87	21.0(0)	442.271 07	442,452 55	442.47613	442.48016	
0	(9)	- 34 D24	440,14849	20.6(1)	442,268 44	442,450 57	442.474.05	442.47808	
	(10)	D_{2h}	440.133.66	18.0(2)	442,253 56	442.441 96	442,464 51	442,467 30	
	(11)	D_{4h}	440.111 41	22.5 (3)	442.229 47	442.42621	442.447 80		

^a Zero-point vibrational energy, unscaled, in kJ mol⁻¹, at HF/3-21G(*). Numbers of imaginary vibrational frequencies are given in parentheses for each stationary point.

Table 2. Relative energies and	binding energies of PLi _n a	and SLi_n molecules (kJ mol ⁻¹).
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		3-21G(*) //3-21G(*)	6-31G(*)//3-21G(*)					
Species			RHF/	MP2/		MP4DQ/	MP4SDQ/	
PL ₁₅	(1)	0.0	0.0	0.0	0.0	0.0		
-	(2)	3.5	3.8	1.1	2.0	2.9		
Δ_1^{a}	()	-175.0	-153.5	-201.6	-190.4	-176.9		
SĹi₄	(3)	0.0	0.0	5.8	3.8	1.2	1.5	
	(4)	3.4	3.8	0.0	0.0	0.0	0.0	
	(5)	6.8	7.5	2.0	2.4	2.8	3.0	
	(6)	10.7	15.4	3.5	8.0	12.2	12.7	
	(7)	43.8	43.8	23.4	25.0			
Δ_2^{b}	~ /	-146.7	-126.4	-173.3	-170.1	-163.1	-164.0	
SLi ₆	(8)	0.0	0.0	0.0	0.0	0.0		
0	(9)	8.9	6.9	5.2	5.5	5.5		
	(10)	47.8	46.0	27.8	30.5	33.8		
	àń	106.2	109.2	69.2	74.4			
Δ_3^{c}	< - /	-51.7	-41.8	-56.0	-54.4	-63.9		

^a Energy of PLi_5 compared to sum of ($PLi_3 + Li_2$). ^b Energy of SLi_4 compared to sum of ($SLi_2 + Li_2$). ^c Energy of SLi_6 compared to sum of ($SLi_4 + Li_2$).

bipyramidal structure, apart from a class containing bidentate chelating ligands with internal ring systems.⁶ Indeed, of all MX₅ species, where M is a main group element and X monodentate, only SbPh₅⁷ and InCl₅²⁻⁸ are known to have square pyramidal structures, and 'packing factors' may well be influential for these solid-state cases, since a trigonal bipyramidal structure is predicted⁹ by *ab initio* methods for AlF₅²⁻.

The occupied valence orbitals for (1) are $6a_1^{2}3e^47a_1^{2}8a_1^2$, which may be contrasted with $4a_1^{2}2e^45a_1^{2}1b_2^{2}$ in PH₅ if $C_{4\nu}$ symmetry is imposed. Orbital $8a_1$ in (1) is bonding between adjacent lithium atoms, but P–Li antibonding. Another $C_{4\nu}$ stationary point for PLi₅ was found inadvertently, corresponding to the $a_1^2 \rightarrow b_1^2$ HOMO \rightarrow LUMO excitation. This second $C_{4\nu}$ state is a transition state, and lies 233 kJ mol⁻¹ above (1) at the 3-21G(*)SCF level. PLi₅ is predicted to be substantially more stable than (PLi₃ + Li₂), as reported in Table 2. The atomisation energy of PLi₃ is calculated to be 416 kJ mol⁻¹ [MP4DQ/6-31G(*)//3-21G(*)], but since solid phosphorus and metallic lithium lie some 820 kJ mol⁻¹ below (P_(g) + 3 Li_(g)), it is probable that PLi₅ is unstable relative to solid phosphorus and metallic lithium. However, its preparation is not necessarily precluded by thermodynamic instability.

No fewer than five different stationary points were located for SLi_4 . Only one of these is a true minimum at 3-21G(*). Interestingly, this isomer (3) has $C_{3\nu}$ symmetry. I believe that there is no precedent for an MX₄ species to have a $C_{3\nu}$ ground state. It is remarkable that three other isomers of SLi_4 lie within 11 kJ mol⁻¹ of (3) at 3-21G(*), especially given the enormous changes in molecular geometry evident in Figure 1. SLi₄ clearly has exceptional fluxional character. It should be noted that only one stationary point, of T_d symmetry, has been reported for OLi_{4} .² The relative energies of (3), (4), and (5) change as more exact theoretical methods are used (see Table 2). Isomer (4) is apparently stabilized over (3) by correlation effects, though since geometry optimisation was not undertaken with correlated wavefunctions this conclusion is tentative for the moment. Although it is not yet possible to predict with confidence which structure is adopted by SLi₄, there can be no doubt that it is more thermodynamically stable than $(SLi_2 + Li_2).$

Four different stationary points were located for SLi₆, of which only one, (8), is a true minimum. Isomer (8) is the most stable isomer of SLi₆ at all levels of theory used here. I know of no other MX₆ molecule with a D_{3d} ground state, with the

probable exception of OLi_6^{10} for which vibrational frequencies have apparently not been calculated. The single D_{3h} imaginary vibrational frequency of 66i cm⁻¹ for (9) is an a" mode which rotates one 'eclipsed' Li₃ prismatic face relative to the other, to give the 'staggered' D_{3d} isomer (8).

The non-octahedral structure of SLi₆ may appear surprising, but can be traced to a Jahn–Teller distortion. Octahedral SLi₆ would have the valence orbital configuration $4a_{1g}^{23}t_{1u}^{65}a_{1g}^{24}t_{1u}^{2}$; the 'extra' valence bonding a_{1g} and t_{1u} orbitals which are possible in SLi₆ but not in SH₆ involve Li 2p orbitals. SLi₆ is predicted to have significant thermodynamic stability compared to (SLi₄ + Li₂), but no structure was found for SLi₈ lying below (SLi₆ + Li₂). Inspection of Figure 1 shows that many of the structures reported here may be regarded as containing triangular Li₃⁺ fragments. Note that the closest Li · · · Li separation is 2.803 Å in (8) and 2.972 Å in (3) compared to 3.055 Å in Li₃⁺ (3-21G/SCF calculations).

While the theoretical methods used here work well for 'standard' molecules,⁵ they may be less applicable to unusual cases. Preliminary MCSCF and MRCI calculations have revealed that while PLi₃ and SLi₂ are reasonably well described by a single reference calculation, the 'hyper-lithiated' species have more complex electronic structures with Hartree–Fock reference coefficients of about 0.93 (PLi₅), 0.94 (SLi₄), and 0.89 (SLi₆). A multireference approach will therefore probably be necessary to give a definitive theoretical

description of these hyperlithiated compounds. Work is continuing in this area.

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